

# SnO<sub>2</sub>-rGO nanocomposite as an efficient electron transport layer for stable perovskite solar cells on AZO substrate

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## Abstract

Electron transport layer (ETL) plays an important role in realizing efficient and stable perovskite solar cells (PSCs). There are continuous efforts in developing new types of low cost ETLs with improved conductivity and compatibility with perovskite and the conducting electrode. Here, in order to obtain high efficient and stable PSCs on ZnO:Al (AZO) substrate, reduced graphene oxide (rGO) is incorporated into SnO<sub>2</sub> nanoparticles to form composite ETL. For planar PSC on AZO substrates, SnO<sub>2</sub>-rGO with a low incorporation ratio of 3 wt% rGO significantly enhances the device short circuit current density ( $J_{sc}$ ) and the fill factor when compared to the device with pristine SnO<sub>2</sub> ETL, leading to an overall power conversion efficiency of 16.8% with negligible hysteresis. The effectiveness of the excited charge transfer process of SnO<sub>2</sub>-rGO ETL is revealed by time-resolved photoluminescence decay, and by electrochemical impedance spectrum measurements. Furthermore, the solar cell stability is also enhanced due to the incorporation of rGO in the ETL. This work provides a low cost and effective ETL modification strategy for achieving high performance planar PSCs.

Supplementary material for this article is available [online](#)

Keywords: AZO substrate, electron transport layer, electrochemical impedance spectroscopy, solar cell, graphene, metal oxide

(Some figures may appear in colour only in the online journal)

## 1. Introduction

Perovskite solar cells (PSCs), with a power conversion efficiency (PCE) reaching 23.3% [1], have attracted great attention in recent years. The remarkable achievement of PSC is owing to the excellent material properties of perovskite, such as suitable bandgap, high carrier mobility and a long diffusion length [2–4]. The improvement of interfacial and hole/electron transporting layer also make a significant contribution to this rapid development. TiO<sub>2</sub> is one of the most widely adopted electron transport layers (ETLs). Nonetheless, the

fact that TiO<sub>2</sub> requires a high temperature sintering process makes it unfavorable for practical application. Even worse, the photocatalytic active TiO<sub>2</sub> could lead to photodegradation of perovskite material [5]. ZnO [6–9] and SnO<sub>2</sub> [10–13] are also being studied as alternative low cost ETL candidates. It is reported that ZnO surface usually contains hydroxide and the residual acetate ligands that can also decompose the perovskite film, adding to the device instability [14]. Among many n-type metal oxide, SnO<sub>2</sub> with suitable conduction band level and excellent mobility is found to be a better electron transportation for PSCs [15]. PSCs with SnO<sub>2</sub> ETL has

achieved a certified efficiency of 20% [16]. A recent report yttrium (Y)-doped  $\text{SnO}_2$  NSAs, results exhibit improved performance with much reduced hysteresis [17]. In addition, lithium (Li) and Stibium (Sb) doped  $\text{SnO}_2$  films also have been successfully reported [18, 19].

Here, we report our investigation on reduced graphene oxide (rGO) incorporated  $\text{SnO}_2$  ETL for PSCs on aluminum doped zinc oxide (AZO) substrate. AZO is abundant, environment friendly and low cost transparent conducting electrodes for optoelectronic devices. Nonetheless, PSCs based on AZO substrate is reported with relative low PCE of  $\sim 13\%$  so far [20, 21], much lower than its ITO and FTO counterparts. The  $\text{SnO}_2$ -rGO were prepared through spin-coating  $\text{SnO}_2$  and rGO mixed solution on the AZO substrate and annealing at  $180^\circ\text{C}$  in air. The rGO possesses excellent charge transport properties, and rGO has proven compatibility with metal oxides [22–24]. However, it has been reported previously that the amount of rGO would affect the photocatalytic properties of the  $\text{SnO}_2$ -rGO suspension solution [25]. It is suspected that with the right amount of rGO in the composite layer, the rGO would enhance the electron transport behavior while constrain its photo-degradation effect on the perovskite thin film. In this study, it is found that the addition of the right amount rGO in the  $\text{SnO}_2$  layer effectively reduce electron trap states and passivate the AZO substrate, making the device more stable under the continuous solar illumination compared to pristine  $\text{SnO}_2$ . The device based on the 3 wt% rGO mass loading in  $\text{SnO}_2$  ETL can reach a PCE of 16.87% with reduced hysteresis.

## 2. Experiment

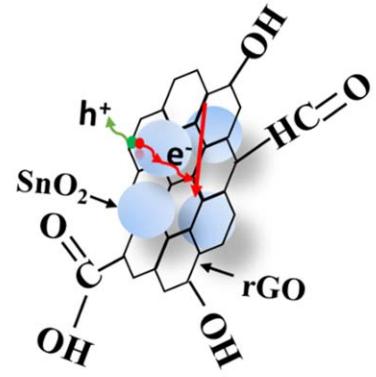
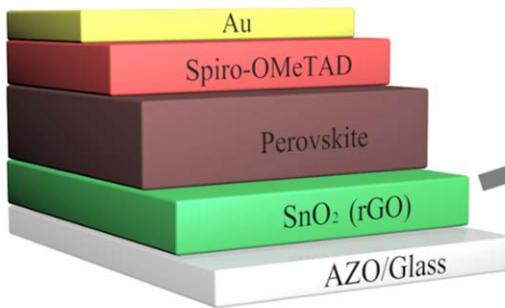
The AZO glass ( $R \approx 10 \Omega \text{ sq}^{-1}$ ,  $T \approx 85\%$ ) was cleaned by deionized water, acetone, ethanol sequentially, and dried by nitrogen purging. Before depositing the  $\text{SnO}_2$  layer, the AZO glass was treated by UV-Ozone for 10 min. The  $\text{SnO}_2$  precursor was purchased from Alfa Aesar (tin oxide, 15% in  $\text{H}_2\text{O}$ ). The  $\text{SnO}_2$  precursor was diluted by  $\text{H}_2\text{O}$  to 3%. The rGO nanosheet (diameter  $\approx 0.5\text{--}5 \mu\text{m}$ , thickness  $\approx 0.8\text{--}1.2 \text{ nm}$ ) was dissolved in GBL with  $1 \text{ mg ml}^{-1}$ , and then mixed with  $\text{SnO}_2$  solution. The prepared suspension solution with different rGO mass percentage (1 wt%, 3 wt%, 5 wt%) were added to the  $\text{SnO}_2$  solution to form our final precursor. The mixed precursor is stable over weeks without any precipitation. The ETL precursor were spin-coated at 3000 rpm for 30 s, followed by annealing at  $180^\circ\text{C}$  for 30 min. The thickness of films were evaluated by using an alpha-step profiler, as shown in figure S1 (available online at [stacks.iop.org/NANO/30/075202/mmedia](https://stacks.iop.org/NANO/30/075202/mmedia)), the average thickness of the resulted ETLs were  $\sim 37 \text{ nm}$ . The results show that the introduction of a small amount of rGO has insignificant influence on the thickness of perovskite films. For the perovskite layer, 159 mg  $\text{CH}_3\text{NH}_3\text{I}$  and 461 mg of  $\text{PbI}_2$  were dissolved in 72 ml of DMSO and 640 ml of DMF solution. The perovskite solution was spin-coated on the  $\text{SnO}_2$  transfer layer at 4000 rpm for 10 s, and then annealed at  $70^\circ\text{C}$  for 10 min and at  $100^\circ\text{C}$  for another 10 min. The HTM solution was preparation by dissolving 72.3 mg

Sprio-OMeTAD in 1 ml chlorobenzene, 18  $\mu\text{l}$  Lithium salt solution ( $520 \text{ mg ml}^{-1}$  dissolved in acetonitrile) and 28  $\mu\text{l}$  4-tert-Butylpyridine were added to Sprio-OMeTAD solution. Finally, 80 nm of gold was deposited through a shadow mask with  $0.1 \text{ cm}^2$  by thermal evaporation. In addition, the PL spectrum were recorded under excitation at 480 nm, and time-resolved photoluminescence (TPRL) collected excitation at 532 nm, measurements were made by using perovskite film deposited on glass/ETLs substrate with room temperature.

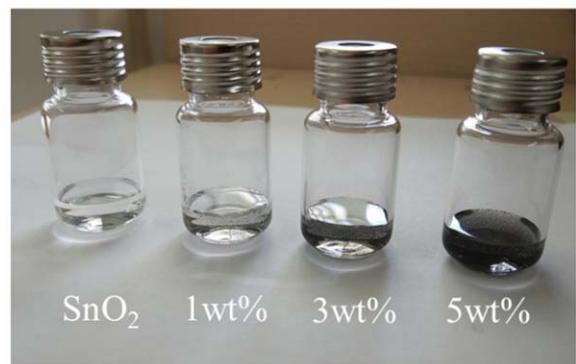
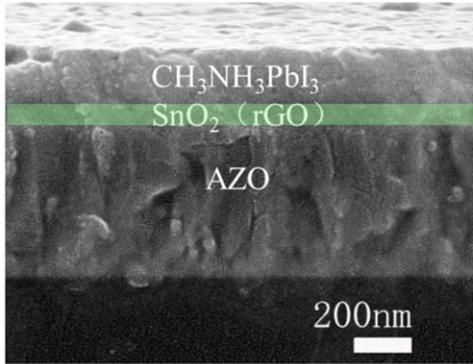
## 3. Results and discussion

Figure 1(a) shows the device structure of AZO/ $\text{SnO}_2$ -rGO/ $\text{CH}_3\text{NH}_3\text{PbI}_3$ /Sprio-OMeTAD/Au, and schematic view of  $\text{SnO}_2$ -rGO mixture. Figure 1(b) is a cross-sectional SEM image illustrates the fabricated device. The perovskite layer consists of the compact microcrystals with the grain size of several hundred nanometers. A uniform  $\text{SnO}_2$ -rGO film of  $\sim 30 \text{ nm}$  forms at the bottom of the AZO substrate. In order to identify the impact of the rGO content on harvesting of the perovskite layer, the optical transmittance of  $\text{SnO}_2$  and  $\text{SnO}_2$ -rGO were measured as a function of rGO content as shown in figure S2. The transmittance decreased with increasing rGO content. It is concluded that excess rGO content would block the light absorption of  $\text{SnO}_2$ , and encourage electron–hole recombination [26]. Furthermore, we monitored the degradation of perovskite layers deposited on  $\text{SnO}_2$ -rGO. Specifically, optical transmittance spectra of perovskite thin films on  $\text{SnO}_2$ -rGO coated substrate were monitored against time in ambient environment with temperature of  $25^\circ\text{C}$  and relative humidity of 50%. The results were presented in the figure S3. Initially, the as deposited perovskite films exhibited black color (figure S3e) and a significant light absorption edge at about 780 nm. It is apparent that the color of the perovskite film on  $\text{SnO}_2$  layer become yellowish after 120 h, revealing very poor stability of the perovskite on  $\text{SnO}_2$  substrate. In contrast, perovskite films with  $\text{SnO}_2$ -rGO (1 wt%) and  $\text{SnO}_2$ -rGO (5 wt%) shows a slower color changing rate. It is worth to note that only slight degradation can be observed for perovskite films deposited on  $\text{SnO}_2$ -rGO (3 wt%) after 120 h (figure S3e). To quantitative the degradation rate, we compared the light transmittance value at the 800 nm ( $T_{800 \text{ nm}}$ ) and the 750 nm ( $T_{750 \text{ nm}}$ ) near the perovskite bandgap absorption edge. The sample with  $\text{SnO}_2$ -rGO (3 wt%) shows the highest  $T_{800 \text{ nm}}/T_{750 \text{ nm}}$  value against time, indicating perovskite thin film on  $\text{SnO}_2$ -rGO (3 wt%) has the best stability under testing condition among these samples (figure S3f). In addition, the average surface roughness of these  $\text{SnO}_2$  films on the AZO substrates were also characterized by atomic force microscope, which is shown in figure S4. The AZO/ $\text{SnO}_2$ -rGO (3 wt%) surface shows the lowest rms roughness ( $\sim 58.1 \text{ nm}$ ), which is lower than the AZO/ $\text{SnO}_2$  substrate surface. To obtain optimal doping concentration, different amount of rGO suspension solution, with different mass percent (0, 1, 3, 5 wt%), were added into  $\text{SnO}_2$  solution (figure 1(c)).

(a)



(b)



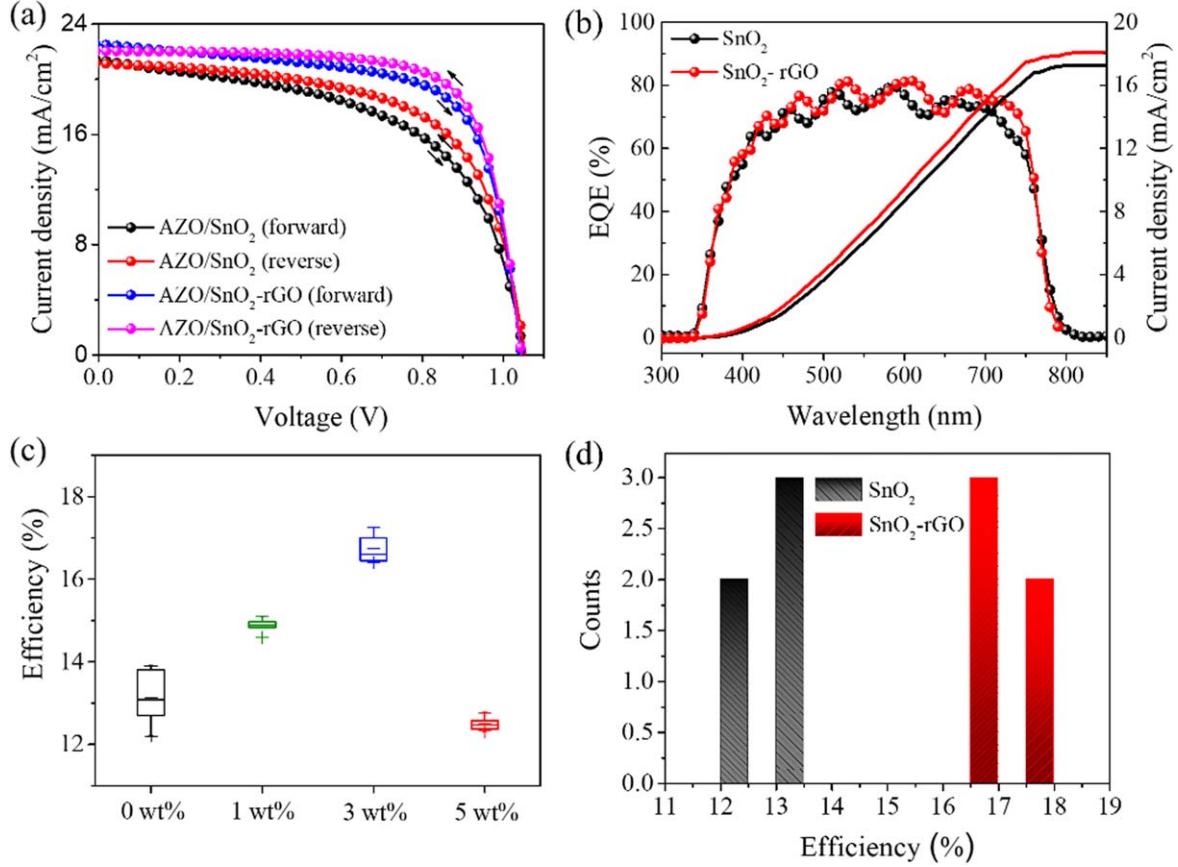
**Figure 1.** (a) Schematic representation of the typical full device structure; (b) cross-sectional SEM image of a complete photovoltaic device based on the AZO/SnO<sub>2</sub>-rGO/CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub> structure (scale bar: 200 nm); (c) the prepared SnO<sub>2</sub>-rGO precursor with rGO mass loading of 0, 1, 3 and 5 wt%.

Figures 2(a) and (b) compares the photocurrent density–voltage (*J*–*V*) characteristics under AM1.5G illumination and the external quantum efficiencies of the champion PSCs with SnO<sub>2</sub>-rGO (3 wt%) and with pristine SnO<sub>2</sub>, respectively. For the champion cells, the *J*<sub>sc</sub> increases from 21.45 to 22.57 mA cm<sup>−2</sup>, and the fill factor increase from 57% to 73%, due to the incorporation of rGO. As a result, the overall PCE is enhanced from 13.9% to 16.87% (table S1). Additionally, in the *J*–*V* characteristic (figure S5), the extracted series resistance (*R*<sub>s</sub>) of the SnO<sub>2</sub>-rGO based device is  $\sim$ 3.78 Ω cm<sup>2</sup>, which is smaller than that of the SnO<sub>2</sub> based device (7.5 Ω cm<sup>2</sup>). The lower *R*<sub>s</sub> of the SnO<sub>2</sub>-rGO ETL can be ascribed to the higher electron mobility than that of the pure SnO<sub>2</sub> ETL. Meanwhile, the shunt resistance (*R*<sub>sh</sub>) for the SnO<sub>2</sub>-rGO device is 2.69 kΩ cm<sup>2</sup>, much larger than that of the SnO<sub>2</sub> device (204 Ω cm<sup>2</sup>), indicating a suppression of the leakage current and reduction of recombination losses in the SnO<sub>2</sub>-rGO device under illumination.

The external quantum efficiency (EQE) spectra of the SnO<sub>2</sub>-rGO and SnO<sub>2</sub> devices are shown in the figure 2(b), presenting EQE values higher than 70% for wavelength between 450 and 750 nm in both cells. Although the integrated *J*<sub>SC</sub> is a bit lower than the results obtain from the *J*–*V* characteristics due to the spectrum mismatch between equipment and theoretical spectra data, the overall *J*<sub>SC</sub> variation trend of

the two devices match well with the *J*–*V* measurement results [27]. The device with SnO<sub>2</sub>-rGO film exhibits higher EQE, compared to the SnO<sub>2</sub> based PSC, leading to a higher integrated *J*<sub>SC</sub>. Since the perovskite films grown on SnO<sub>2</sub>-rGO film and SnO<sub>2</sub> film exhibits UV–vis absorbance (figure S6), it can be inferred that the *J*<sub>sc</sub> enhancement mainly derives from the improved electron transfer efficiency via rGO at the perovskite and ETL interface. It is found that the additive of 3 wt% rGO loading into SnO<sub>2</sub> precursor can significantly improve the performances of the devices. In fact, the device doped with 3 wt% rGO in SnO<sub>2</sub> presents the best results, with an averaged PCE of  $\sim$ 16.7% as shown in the statistic results in figure 2(c). Figure 2(d) is a histogram of PCE over 20 devices fabricated on the SnO<sub>2</sub> and SnO<sub>2</sub>-rGO, obtaining that the devices based on the SnO<sub>2</sub>-rGO an average PCE of  $16.4 \pm 0.4\%$ . Hence, it is deduced that the incorporation of rGO bring a more efficiency and stable PSCs.

To further investigate the charge-transport kinetics at the SnO<sub>2</sub>-rGO/MAPbI<sub>3</sub> interface, the steady-state PL spectra and TRPL measurement are employed. The PL spectra is collected from perovskite films fabricated on glass/ETLs in the figure 3(a), PL peak are observed about 780 nm for both perovskite films with or without rGO. The SnO<sub>2</sub>-rGO/MAPbI<sub>3</sub> film has a much lower PL quantum yield than the SnO<sub>2</sub>/MAPbI<sub>3</sub> film, indicating a higher charge extraction rate across the



**Figure 2.** (a)  $J$ – $V$  curves in forward bias and reverse bias for the best-performing PSC solar cell with the doping concentration of 3 wt% (SnO<sub>2</sub>-rGO) and the reference one (SnO<sub>2</sub>); (b) EQE spectra of the optimal SnO<sub>2</sub> and SnO<sub>2</sub>-rGO based devices; (c) PCE distribution of the devices with and without rGO in the SnO<sub>2</sub> ETL; (d) histograms of PCE with the doping concentration of 3 wt% (SnO<sub>2</sub>-rGO) and the reference SnO<sub>2</sub> ETL.

interface between the SnO<sub>2</sub>-rGO and MAPbI<sub>3</sub>. The TRPL was measured, adopting the same structures as used for the PL measurements. The TRPL results of devices based on SnO<sub>2</sub> and SnO<sub>2</sub>-rGO ETL are presented in figure 3(b). The two TRPL curves can be fitted by a bi-exponential equation [27, 28]:

$$f(t) = A_1 e^{-t/\tau_1} + A_2 e^{-t/\tau_2} + B, \quad (1)$$

where the terms  $A_1$ ,  $A_2$  are the decay amplitude,  $B$  is the constant for the base-line offset,  $\tau$  is the decay time. The decay portion including a fast decay ( $\tau_1$ ) component reflects the quenching of free carriers in the interface between the perovskite and ETL, and a slow decay ( $\tau_2$ ) component represents the recombination of free carriers in radiative channels [28]. The PL decay of the perovskite deposited on the SnO<sub>2</sub> film presents a fast decay lifetime ( $\tau_1$ ) of 31 ns and a slow decay lifetime ( $\tau_2$ ) of 6.4 ns. In contrast, for the SnO<sub>2</sub>-rGO/MAPbI<sub>3</sub> film, the fast decay lifetime and the slow decay lifetime are both shortened, with extracted values of 23 ns and 5.2 ns respectively. The faster PL quenching rate indicates a faster charge transfer rate at the SnO<sub>2</sub>-rGO/MAPbI<sub>3</sub> junction interface. The PL spectra and the TPRL results are consistent with our finding from  $J$ – $V$  and EQE measurements, suggesting that the addition of the rGO in the SnO<sub>2</sub> ETL, the  $J_{sc}$  is improved.

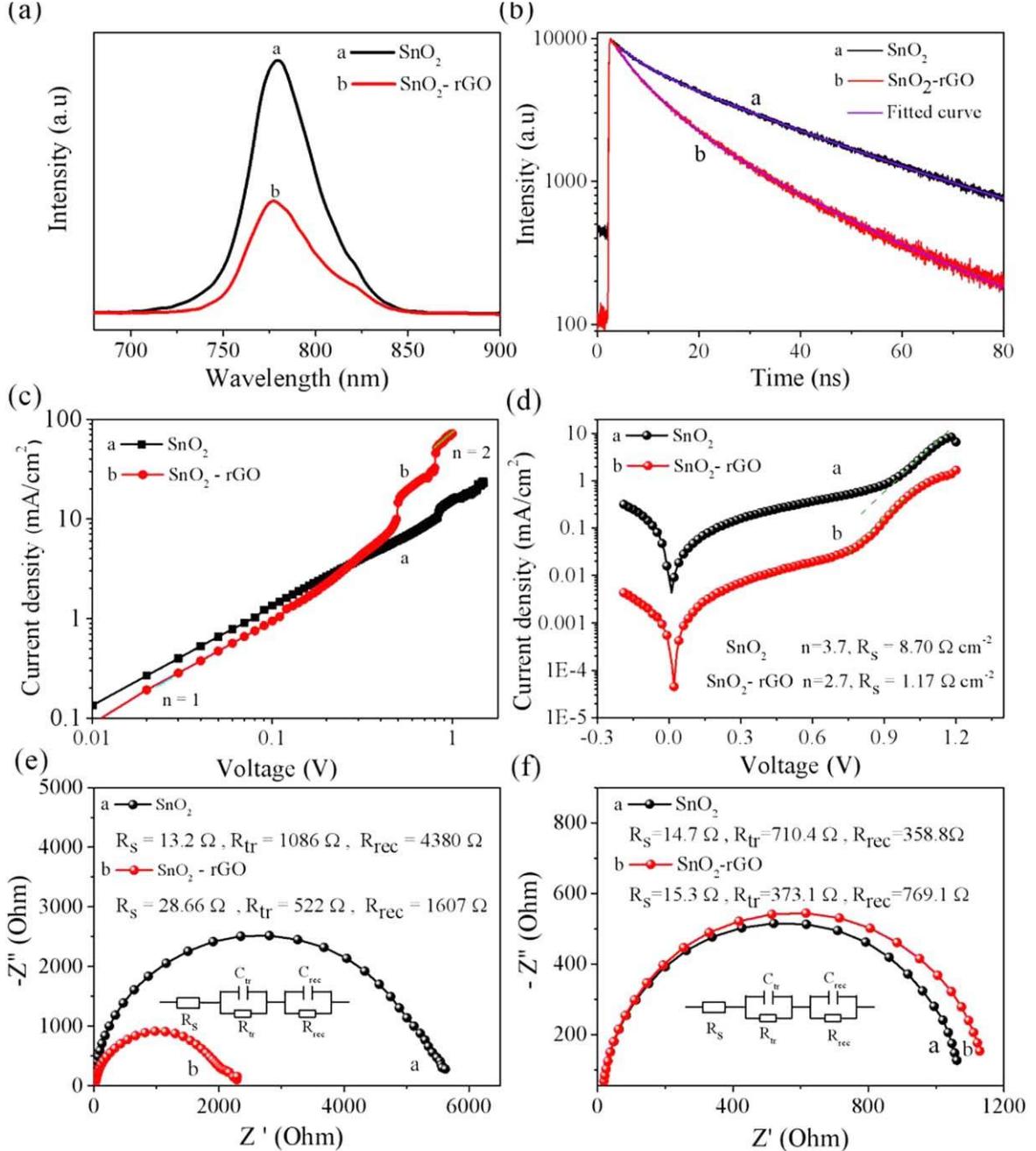
A space-charge limited current (SCLC) measurements were also carried out to evaluate the electron mobility of SnO<sub>2</sub>-rGO

ETL, the results of which are shown in figure 3(c). We have fabricated device with AZO/SnO<sub>2</sub>-rGO/MAPbI<sub>3</sub>/PCBM/BCP/Au configuration for electron transport investigation. The  $J$ – $V$  curves from electron-transport-only measurements can be cataloged into three regions according to the value of the exponent  $n$  [29], i.e.  $n = 1$  is the ohmic region,  $n = 2$  is the SCLC region, and the intermediate region between 1 and 2 is the trap-filled limited region. In the SCLC region, the electron mobility of the PSCs with SnO<sub>2</sub>-rGO film can be evaluated on the basis of the Mott–Gurney law [29]

$$J = \frac{9}{8} \varepsilon_0 \varepsilon \mu \frac{V^2}{d^3}, \quad (2)$$

where  $V$  is the applied voltage,  $d$  is the thickness of the ETLs,  $\varepsilon$  is the dielectric constant of the SnO<sub>2</sub> (12.5) film and  $\mu$  is the electron mobility. The estimated electron mobility of  $3.3 \times 10^{-3} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  can be obtained from the equation (2) from the device based SnO<sub>2</sub>-rGO film, which is about 5 times higher than that of the SnO<sub>2</sub> based device ( $7 \times 10^{-4} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ ). The magnitudes of our mobility values are in the same order to those reported in reference [30]. The SCLC results further confirms that the photogenerated charge carriers are more efficiently transferred via the SnO<sub>2</sub>-rGO ETL than the pristine SnO<sub>2</sub> ETL.

Dark current measure is a method to evaluate charge carrier loss through leakage pathways and the recombination of free



**Figure 3.** (a) Photoluminescence emission spectra of AZO/SnO<sub>2</sub>/perovskite and AZO/SnO<sub>2</sub>-rGO/perovskite; (b) TRPL of CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub> contacted with different interfaces: SnO<sub>2</sub>, SnO<sub>2</sub>-rGO; (c) *I*-*V* characteristics of the AZO/SnO<sub>2</sub>/CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub>/PCBM/BCP/Au and AZO/SnO<sub>2</sub>-rGO/CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub>/PCBM/BCP/Au; (d) *J*-*V* curves for the best-performing PSC solar cell with the doping concentration of 3 wt% (SnO<sub>2</sub>-rGO) and the reference one (SnO<sub>2</sub>) in the dark. Electrochemical impedance spectroscopy of PSC based on SnO<sub>2</sub> and SnO<sub>2</sub>-rGO films under the dark (e) and illumination condition(f).

carriers in devices. Figure 3(d) shows the dark *I*-*V* curves for the best PSCs with SnO<sub>2</sub> and with SnO<sub>2</sub>-rGO, the *J*-*V* characteristics in dark can be described by the Shockley diode equation [31]:

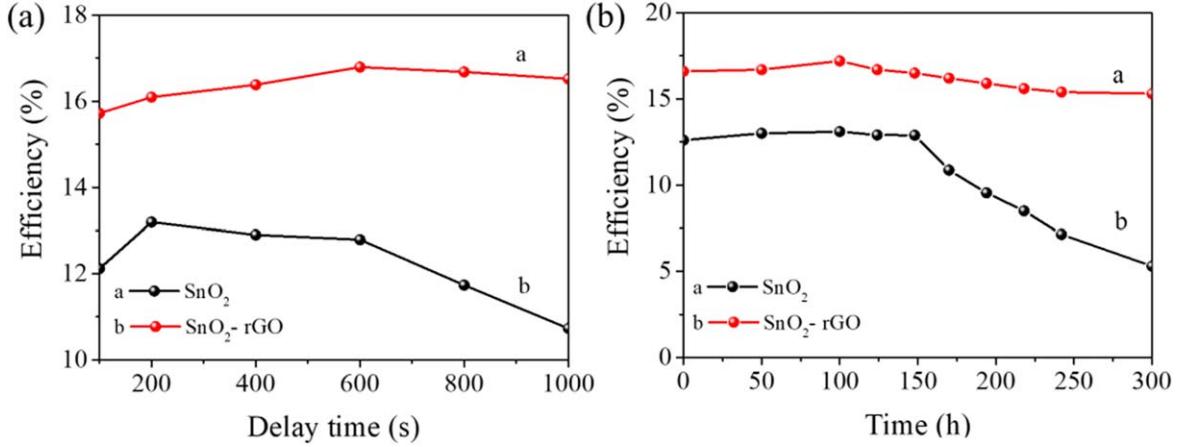
$$J = J_0 \left[ \exp \left( \frac{qV}{nKT} \right) - 1 \right] \quad (3)$$

where  $J_0$  is the saturation current density,  $V$  is the applied voltage,  $n$  is the ideality factor,  $T$  is the temperature and  $K$  is the Boltzmann constant. The dark curves show typical diode properties,

with a transition to the SCLR at the build in voltage ( $V_{bi}$ ). The ideality factor can be determined from the slope of the exponential regime on the semi-logarithmic plot and can be denoted as:

$$n = \left( \frac{kT}{q} \frac{\partial \ln J}{\partial V} \right)^{-1}. \quad (4)$$

In the equation (4),  $n$  of the device with SnO<sub>2</sub> has a value of 3.7, which is higher than that of the device with SnO<sub>2</sub>-rGO



**Figure 4.** (a) SnO<sub>2</sub> (curve a) and SnO<sub>2</sub>-rGO (curve b) ETLs based PSCs under continuous solar illumination; (b) devices stability with SnO<sub>2</sub> and SnO<sub>2</sub>-rGO ETL in the glove box without encapsulation.

of 2.7.  $n$  extracted value are significantly large than 1 for both devices, demonstrating that trap assisted recombination plays a mass of role in the recombination process [30]. However, PSC based SnO<sub>2</sub>-rGO ETL showed an  $n$  value and a  $R_s$  ( $1.17 \Omega \text{ cm}^2$ ) smaller with regard to the SnO<sub>2</sub> ETL ( $8.7 \Omega \text{ cm}^2$ ). Suggesting a decreased trap state density at AZO/SnO<sub>2</sub> interface when rGO is added to SnO<sub>2</sub>. According to equation (3), it can be obtained that the  $J_0$  of the device without rGO is  $6.32 \times 10^{-5} \text{ mA cm}^{-2}$ , which is one order of magnitude large than that of the device with SnO<sub>2</sub>-rGO is  $1.3 \times 10^{-5} \text{ mA cm}^{-2}$ . This decrease in  $J_0$  indicates that carrier recombination is suppressed by the use of rGO. It is suspected that the rGO helps in forming a more compact ETL by filling the voids that originated from SnO<sub>2</sub> nanoparticle aggregations, thus reducing the leakage paths by preventing the perovskite layer from direct contacting the AZO surfaces.

To further reveal the role of the rGO in SnO<sub>2</sub> film, electrochemical impedance spectroscopy analysis was applied to the devices with SnO<sub>2</sub> and SnO<sub>2</sub>-rGO ETL under dark and illumination conditions with a bias voltage of 0.1 V. The Nyquist curves of both devices under dark condition have been plotted in figure 3(e), with an equivalent circuit shown as inset diagram. The resistance  $R_s$  can be extracted from the intersections of these arcs, corresponding to the resistance of AZO/ETLs and HTLs/Au. The resistance  $R_{tr}$  from the high frequency arc regime mostly attributed to the transport resistance of ETLs/Perovskite/Spiro-OMeTAD layers [32, 33]. The fitted value of the ohmic resistance ( $R_s$ ) and charge transfer resistance ( $R_{tr}$ ) in the figure 3(e). The  $R_{tr}$  of the device based on SnO<sub>2</sub> PSCs is  $1086 \Omega$ , and the  $R_{tr}$  of the device based on SnO<sub>2</sub>-rGO is  $522 \Omega$ . The device based SnO<sub>2</sub>-rGO exhibits the smaller charge transfer resistance suggests that electron transport faster at the SnO<sub>2</sub>-rGO/Perovskite interface. Figure 3(f) shows the Nyquist plots of both SnO<sub>2</sub> and SnO<sub>2</sub>-rGO based device under illumination condition. The resistance from lower frequency arc is related to the recombination resistance ( $R_{rec}$ ). The  $R_{rec}$  of the SnO<sub>2</sub>-rGO device is  $769.1 \Omega$ , and the  $R_{rec}$  of the SnO<sub>2</sub> device is  $358.8 \Omega$ . The calculated values of the charge transfer resistance ( $R_{tr}$ ) using the inset equivalent circuit is

$373.1 \Omega$  for the device based on SnO<sub>2</sub>-rGO, which is much lower than the device based on SnO<sub>2</sub> with  $710.4 \Omega$ . The carrier diffusion length  $L_d$  can be estimated by [34]

$$L_d = \left( \frac{R_{rec}}{R_{tr}} \right)^{\frac{1}{2}} L, \quad (5)$$

where  $R_{rec}$  is the recombination resistance,  $R_{tr}$  is the charge transfer resistance,  $L$  is the layer thickness. According to the equation (5), with similar  $L$  in the devices, the large  $R_{rec}$  and small  $R_{tr}$  in solar cell decrease the possibility for electron–hole recombination.

In order to confirm light stability, the light stability measurements on the PSC were investigated. The  $J$ – $V$  measure of PSC with SnO<sub>2</sub>-rGO electron transfer layer under different continuous illumination time from 100 to 1000 s were measured and presented in the figure 4(a). The SnO<sub>2</sub> based PSC degrades with time rapidly, with PCE decreasing from  $\sim 12\%$  to  $11\%$  within 1000 s. The most prominent decay is obtained in the plot of  $V_{oc}$  (in the figure S7). In contrast, upon light exposure, the SnO<sub>2</sub>-rGO device remains more stable with PCE above 16% for more than 1000 s. The stability of devices was also tested as a function of storage time in the glove box without encapsulation. The PCE is plotted independently in the figure 4(b) to reveal the stability of the devices clearly. The SnO<sub>2</sub>-rGO device, the PCE keep 80% after 300 h while PCE of the SnO<sub>2</sub> device drops down to  $<50\%$ . It is worth noting that the stability of PSCs was tested under AM. 1.5G as light source. The results confirm that the SnO<sub>2</sub>-rGO based PSCs is more stable than the SnO<sub>2</sub> device.

The enhanced stability of PSC may be ascribed to the passivation effect of rGO at the interface. For example, the surface of the AZO and low temperature solution processed SnO<sub>2</sub> contain many trap states originated from oxygen vacancies [35], which leads to decomposed of perovskite at the interface and increase carrier recombination. The rGO is suspected to distribute at the boundary of SnO<sub>2</sub> nanocrystals and passivate the electron traps of the SnO<sub>2</sub> and the AZO surface.

## 4. Conclusion

In this study, we have synthesized a new way to improve the performance of planar n-i-p PSCs by adding rGO of 3 wt% into  $\text{SnO}_2$ . The addition of rGO resulted in a high efficiency of 16.87% with a very reduced hysteresis, and the devices deposited on the AZO substrate and treated below 180 °C. In addition, the increased  $J_{sc}$  of the device based on  $\text{SnO}_2$ -rGO, which attributes to effective electron extraction from perovskite to AZO. On the other hand, the rGO doped  $\text{SnO}_2$  can effectively improve the device stability is due to the stabilization of AZO/ $\text{SnO}_2$  interface. The introduction of rGO in  $\text{SnO}_2$  is a simple and effective way to improve device performances.

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